LASER INTERFEROMETER GRAVITATIONAL WAVE OBSERVATORY
- LIGO -
CALIFORNIA INSTITUTE OF TECHNOLOGY
MASSACHUSETTS INSTITUTE OF TECHNOLOGY

Document Type	**LIGO-T010110-00-Z**	2014/04/09

Beam Pattern Response Functions and Times of Arrival for Earthbound Interferometer

Warren Anderson, Patrick Brady, David Chin, Jolien Creighton, Keith Riles, and John Whelan

Distribution of this draft:

All

California Institute of Technology
LIGO Project - MS 51-33
Pasadena CA 91125
Phone (626) 395-2129
Fax (626) 304-9834
E-mail: info@ligo.caltech.edu

Massachusetts Institute of Technology
LIGO Project - MS NW22-295
Cambridge, MA 01239
Phone (617) 253-4824
Fax (617) 253-7014
E-mail: info@ligo.mit.edu

WWW: http://www.ligo.caltech.edu/

Processed with LaTeX on 2014/04/09

Beam Pattern Response Functions and Times of Arrival for Earthbound Interferometers

Warren G. Anderson and John T. Whelan (The University of Texas at Brownsville),
Patrick R. Brady and Jolien D. E. Creighton (University of Wisconsin - Milwaukee),
David Chin and Keith Riles (University of Michigan)

First Version March 30, 2000; Last Revised July 18, 2002

Preliminaries

Notes on usage

Hopefully everything one needs in order to calculate response functions and times of arrival is contained in this worksheet, along with intermediate results which one can use to check calculations. There are a few results in this worksheet that are so long and unweildy that we chose to supress the output by using a colon (:) to end the input statement rather than the usual semicolon (;). We always warn you that this is the case by putting (results supressed) into the preceding text. If you need to see these results, just change the colon to a semicolon.

Technical asides

LIGO technical document LIGO-T980044-08-E

Much of this document is based upon LIGO technical document LIGO-T980044-08-E. There are, however, several definitions that are left to the reader to determine, presumably because their is only one clear choice. We have made those definitions obvious here, to wit:
- the LIGO document refers to a suface, a point, and a line through the point prependiular to the surface. There can in principle be an arbitrary number of such lines. Clearly, they mean the line passing through the surface whose intersection with the surface is closest to the point. It can be shown that this line always intersects the surface perpendicularly.
- the term local horizontal at a point is undefined. We have taken this to mean the plane through that point which is parallel to the tangent to the surface at the intersection point mentioned above.
- they give values for angles of vectors "above the local horizontal". Again, there are an infinite number of such angles. Clearly, they mean the angle of smallest magnitude.

We have also corrected one typographical error in the document - the denominator of the equation for $R[\phi]$ should be the square root of the value quoted.

Finally, the document also refers to an eccentricity, ε. However, it is a quantity derived from two other quoted quantities, a and b ($[\ 1-\varepsilon^2\ ]=\frac{b^2}{a^2}$), and is therefore unnecessary, so we have not used it.

Note about GMST

In the Times of Arrival section, one encounters the following situation: the vector describing the location of the detector depends on the GMST at which the wave arrives at the detector. However, the whole purpose of this section is to calculate the difference between the GMST at which the wave arrives at the detector and the GMST at which it arrives at the center of the earth. This could be dealt with by inverting the functions of GMST, at the cost of computational complexity. However, the dependence of the time of arrival on this difference between the two GMST's is clearly a higher order effect, and we will ignore it. Note however that this may not be the case for the response functions themselves. The earth can rotate by approximately 0.3 arcseconds during the time it takes light to travel from the center of the earth to it's limb. This might have a significant effect on the response functions for some applications.

Housekeeping duties

Getting started. restart for a clean slate. Load the linalg and tensor packages to give us matrix and tensor handling capabilities.

```
> restart; with(linalg):with(tensor):
Warning, new definition for norm
Warning, new definition for trace
```

Earth Model WGS-84

It is convenient to first specify an earth model with which to relate the detector position and orientation to the gravitational wave propagation direction and polarization. The earth model used here is the same one used in LIGO technical document LIGO-T980044-08-E; earth model WGS-84. The reference figure of this model is the surface of (Σ) an oblate ellipsoid with semi-major axis $a=6378137\ m$, semi-minor axis $b=.6356752314\ 10^{7}\ m$. The document also refers to the eccentricity, ε. However, it is a quantity derived from a and b (i.e. $[\ 1-\varepsilon^2\ ]=\frac{b^2}{a^2}$), and it is therefore unnecessary. The set of relevant parameters is therefore (in units of meters):

```
> WGS84_params:={a = 6378137,b = 6356752.314};
```

$$WGS84\_params := \{ a=6378137,\ b=.6356752314\ 10^{7} \}$$

In WGS-84, one specifies the position of any point in space x in terms of the point $\sigma(\ x\ )$, which is defined to be the closest point on the surface Σ to x. The coordinates of x are then specified by l and λ, the North latitude and East longitude of $\sigma(\ x\ )$ respectively, and by the height h, which is the distance from $\sigma(\ x\ )$ to x along the outward normal to Σ. *N.B. the line segment connecting any point in space to the nearest point on any smooth surface intersects the surface orthogonally.*

Coordinate Frames

There are at least four coordinate frames which are useful in discussing the detection of gravitational waves:

- source frame, which is chosen to provide a simple description of the source
- the wave propagation frame, which provides a simple description of the waves propagating from the source to the earth
- the earth fixed frame, which provides a standard earthbound coordinate frame in which to describe the waves
- the detector frame, which is adapted to the particular geometry of the detector which is to detect the waves

This worksheet is primarily concerned with the latter three frames. Each frame is an orthogonal Cartesian frame. It is assumed that the spatial curvature of space-time near the earth is sufficiently small that Euclidean translations may be used to relate frames at the earth’s center and its surface.

Wave Propagation Frame

The wave propagation frame is chosen to be compatible with the definition of Will and Wiseman. The origin of this frame is taken to be the centroid of the earth. The z-axis points lies along the line joining the origin to the source, and points *away* from the source. The x-axis will typically be chosen to make the transition from the source frame to the wave propagation frame convenient. For the purposes of this document, the x-axis can be considered to be an arbitrary vector chosen in the plane orthogonal to the z-axis. The y-axis is then chosen to complete a right-handed coordinate system. The coordinates in the wave propagation frame are denoted by (x_w, y_w, z_w). The unit vectors are:

```
> e_wx:=create([-1],array(1..3,[1,0,0]));e_wy:=create([-1],array(1..
  3,[0,1,0]));e_wz:=create([-1],array(1..3,[0,0,1]));
```

```
e_wx := table([
    index_char = [ -1 ]
    compts = [ 1, 0, 0 ]
    ])
e_wy := table([
    index_char = [ -1 ]
    compts = [ 0, 1, 0 ]
    ])
e_wz := table([
    index_char = [ -1 ]
    compts = [ 0, 0, 1 ]
    ])
```

Earth Fixed Frame

The earth fixed frame is defined in LIGO technical document LIGO-T980044-08-E. The origin

of this frame is the centroid of the earth. The z-axis points from the origin to the North pole $\{ l = 90\ N,\ \lambda = 0 \}$. The x-axis points from the origin to the intersection of the earth's equator and prime meridian $\{ 0,\ 0 \}$. The y-axis is chosen to complete a right-handed coordinate system $\{ 0,\ 90\ E \}$. Note that this coordinate system rotates with earth (with respect to the fixed background stars). Its relationship to the wave propagation frame therefore changes as a periodic function of time, with a period of one sidereal day. The coordinates in the earth fixed frame are denoted by
($x_e,\ y_e,\ z_e$).The unit vectors are:

```
> e_ex:=create([-1],array(1..3,[1,0,0]));e_ey:=create([-1],array(1..
  3,[0,1,0]));e_ez:=create([-1],array(1..3,[0,0,1]));
```

```
e_ex := table([
    index_char = [ -1 ]
    compts = [ 1, 0, 0 ]
    ])
e_ey := table([
    index_char = [ -1 ]
    compts = [ 0, 1, 0 ]
    ])
e_ez := table([
    index_char = [ -1 ]
    compts = [ 0, 0, 1 ]
    ])
```

Detector Frame

We define the detector frame in terms of cardinal compass points at the position of the detector on the earth's surface. The definition of the detector frame assumes the oblate ellipsoidal earth model WGS-84. We begin with some preliminary definitions. Let $\sigma(x)$ be the on the reference ellipsoid's surface (Σ) nearest to the point x. Let $\Lambda(x)$ be the tangent plane to Σ at $\sigma(x)$. Define the *local horizontal* at x to be the plane containing x which is parallel to $\Lambda(x)$. Within the local horizontal, the directions North and East are inherited from their definitions in $\Lambda(x)$. Note that there is no ambiguity in this since the local horizontal is parallel to $\Lambda(x)$. $\sigma(x)$ is the point on the surface of the ellipsoid that is closest to the corner of the interferometer (this will be unique if the corner is near enough the surface). The origin of the detector frame is then the corner of the interferometer. The $x-y$ plane is defined to be the local horizontal. The x-axis is chosen to point due East in this plane, and the y-axis to point due North, so that x and y are orthogonal. The z-axis points along the outward normal to the earth's surface. The coordinates in the detector frame are denoted by ($x_d,\ y_d,\ z_d$).The unit vectors are:

```
> e_dx:=create([-1],array(1..3,[1,0,0]));e_dy:=create([-1],array(1..
```

```
3,[0,1,0]));e_dz:=create([-1],array(1..3,[0,0,1]));
e_dx := table([
   index_char = [ -1 ]
   compts = [ 1, 0, 0 ]
   ])
e_dy := table([
   index_char = [ -1 ]
   compts = [ 0, 1, 0 ]
   ])
e_dz := table([
   index_char = [ -1 ]
   compts = [ 0, 0, 1 ]
   ])
```

Gravitational Wave Tensor Components

The components of the wave are naturally given in the wave propagation frame. As a first step, define the polarization tensors `e_plus` and `e_cross` to be:

```
> e_plus:=lin_com(1,prod(e_wx,e_wx),-1,prod(e_wy,e_wy));
e_plus := table([
   index_char = [ -1, -1 ]
             [1    0    0]
   compts =  [0   -1    0]
             [0    0    0]
   ])
```

```
> e_cross:=lin_com(1,prod(e_wx,e_wy),1,prod(e_wy,e_wx));
e_cross := table([
   index_char = [ -1, -1 ]
             [0   1   0]
   compts =  [1   0   0]
             [0   0   0]
   ])
```

The gravitational wave tensor is a linear combination of e_{plus} and e_{cross}

```
> h_tensor:=lin_com(h_plus,e_plus,h_cross,e_cross);
h_tensor := table([
   index_char = [ -1, -1 ]
```

$$compts = \begin{bmatrix} h\_plus & h\_cross & 0 \\ h\_cross & -h\_plus & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

```
])
```

The Transformation from Wave Propagation to Earth Fixed Frames

We will need to write an arbitrary vector in the wave propagation frame in terms of earth fixed frame unit vectors. The transformation between the wave propagation and earth fixed frames that will allow us to do this is carried out using Euler angle rotations.

Definition of the Euler Angles

First, some initial definitions. "Counterclockwise about an axis" means in a counterclockwise direction as viewed from that axis while facing the origin. The line of nodes is defined to be the line of intersection of the $x_e - y_e$ plane and the $x_w - y_w$ plane. If one draws a circle about the origin on the $x_e - y_e$ plane, their will be two points of intersection between the circle and the line of nodes, and these are called nodes. Draw a circle in the $x_w - y_w$ plane about the origin and passing through the nodes. Draw tangent vectors to the circle pointing counterclockwise about the z_w-axis. At one node, the tangent vector thus drawn will have a positive z_e component, at the other the tangent vector will have a negative z_e component. The former is called the ascending node, and the half line from the origin through the ascending node is called the line of ascending nodes. The Euler angles are defined with respect to the axes of the two frames and the line of ascending nodes. They are:

Ψ - the angle from the line of ascending nodes to the positive x_w-axis counterclockwise about the positive z_w-axis.

Θ - the smallest angle between the positive z_e-axis and the positive z_w-axis. It is always positive.

Φ - the angle from the positive x_e-axis to the line of ascending nodes counterclockwise about the z_e-axis.

The Rotation Matrices

The three rotation matrices are:

```
> R[Psi]:=array(1..3,1..3,[[cos(Psi),sin(Psi),0],[-sin(Psi),cos(Psi)
,0],[0,0,1]]);
```

$$R_{\Psi} := \begin{bmatrix} \cos(\Psi) & \sin(\Psi) & 0 \\ -\sin(\Psi) & \cos(\Psi) & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

```
> R[Theta]:=array(1..3,1..3,[[1,0,0],[0,cos(Theta),sin(Theta)],[0,-s
in(Theta),cos(Theta)]]);
```

$$R_{\Theta} := \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos(\Theta) & \sin(\Theta) \\ 0 & -\sin(\Theta) & \cos(\Theta) \end{bmatrix}$$

```
> R[Phi]:=array(1..3,1..3,[[cos(Phi),sin(Phi),0],[-sin(Phi),cos(Phi)
,0],[0,0,1]]);
```

$$R_{\Phi} := \begin{bmatrix} \cos(\Phi) & \sin(\Phi) & 0 \\ -\sin(\Phi) & \cos(\Phi) & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

The full rotation matrix which transforms wave propagation frame vectors in the fixed earth frame is:

```
> R[Euler]:=multiply(R[Psi],multiply(R[Theta],R[Phi]));
```

$$R_{Euler} :=$$

$$\begin{aligned} &[\cos(\Psi)\cos(\Phi) - \sin(\Psi)\cos(\Theta)\sin(\Phi)\,, \\ &\cos(\Psi)\sin(\Phi) + \sin(\Psi)\cos(\Theta)\cos(\Phi)\,,\ \sin(\Psi)\sin(\Theta)] \\ &[-\sin(\Psi)\cos(\Phi) - \cos(\Psi)\cos(\Theta)\sin(\Phi)\,, \\ &-\sin(\Psi)\sin(\Phi) + \cos(\Psi)\cos(\Theta)\cos(\Phi)\,,\ \cos(\Psi)\sin(\Theta)] \\ &[\sin(\Theta)\sin(\Phi)\,,\ -\sin(\Theta)\cos(\Phi)\,,\ \cos(\Theta)] \end{aligned}$$

Sky Angles

While the Euler angles are useful for defining the rotation matrix, the gravitational wave community has historically used angles specifying the sky position rather than the Euler angles above. We will recalculate in terms of these angles now.

Definition of Sky Angles

The three angles in terms of which the response functions are usually calculated are the altitude θ, the azimuth ϕ, and the polarization angle ψ. Altitude and azimuth specify the position of the gravitational wave source on the sky with respect to the earth fixed coordinates, i.e. they define a vector in the earth fixed frame that points toward the source. The polarization angle ψ contains information about the polarization of the wave. Altitude and azimuth might best be understood by their simple relationship to declination δ and right ascension α. The relationship is

```
> sky_to_celestial:={phi=alpha-GMST,theta=Pi/2-delta};
```

$$sky\_to\_celestial := \{\phi = \alpha - GMST,\ \theta = \frac{1}{2}\pi - \delta\}$$

where GMST, the Greenwich Mean Sidereal Time of the observation, allows us to relate a frame fixed with respect to the celestial sphere to the rotating earth fixed frame.

Relationship Between Sky Angles and Euler Angles

The relationship between the sky angles and the Euler angles is also particularly simple. Recall that Θ measures the angle between the z_e-axis and the z_w-axis, which points *away* from the source. Thus, the angle θ between the z_e and the direction *to* the source must be given by $\theta=\pi-\Theta$. Likewise, recall that Φ is the angle from the x_e-axis to the line of ascending nodes. Furthermore, observe that the projection of the z_w-axis onto the x_e-y_e plane is $\frac{\pi}{2}$ clockwise from the line of ascending nodes about the positive z_e-axis. Thus, ϕ, the angle from the x_e-axis to the line of ascending nodes about the positive z_e-axis, is given by $\phi=\Phi+\frac{\pi}{2}$. Finally, since the Euler angle Ψ encodes all the necessary polarization information already, we take it to be the polarization angle, i.e. $\psi=\Psi$. The complete coordinate transformation is therefore

```
> Euler_to_sky:={Theta=Pi-theta, Phi=phi-Pi/2, Psi=psi};
```

$$Euler\_to\_sky := \{\Theta=\pi-\theta,\ \Phi=\phi-\frac{1}{2}\pi,\ \Psi=\psi\}$$

The Rotation Matrices in Terms of Sky Angles

It is now a straightforward matter to express the rotation matrices in terms of the sky angles. They are:

```
> R[psi]:=map2(subs,Euler_to_sky,R[Psi]);
```

$$R_\psi := \begin{bmatrix} \cos(\psi) & \sin(\psi) & 0 \\ -\sin(\psi) & \cos(\psi) & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

```
> R[theta]:=map(expand,map2(subs,Euler_to_sky,R[Theta]));
```

$$R_\theta := \begin{bmatrix} 1 & 0 & 0 \\ 0 & -\cos(\theta) & \sin(\theta) \\ 0 & -\sin(\theta) & -\cos(\theta) \end{bmatrix}$$

```
> R[phi]:=map(expand,map2(subs,Euler_to_sky,R[Phi]));
```

$$R_\phi := \begin{bmatrix} \sin(\phi) & -\cos(\phi) & 0 \\ \cos(\phi) & \sin(\phi) & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

```
> R[sky]:=map(expand,map2(subs,Euler_to_sky,R[Euler]));
```

$$R_{sky} :=$$

$$[\cos(\psi)\sin(\phi)-\sin(\psi)\cos(\theta)\cos(\phi),$$
$$-\cos(\psi)\cos(\phi)-\sin(\psi)\sin(\phi)\cos(\theta),\ \sin(\psi)\sin(\theta)]$$
$$[-\sin(\psi)\sin(\phi)-\cos(\psi)\cos(\theta)\cos(\phi),$$
$$\sin(\psi)\cos(\phi)-\cos(\psi)\sin(\phi)\cos(\theta),\ \cos(\psi)\sin(\theta)]$$
$$[-\sin(\theta)\cos(\phi),\ -\sin(\theta)\sin(\phi),\ -\cos(\theta)]$$

The Gravitational Wave Tensor in the Earth Fixed Frame

In order to find detector response functions, it is convenient to express the gravitational wave vector in terms of the earth fixed frame coordinate basis. To do this, we first create a triad to hold the earth fixed frame vectors

```
> triad[e]:=vector([e_e_x,e_e_y,e_e_z]);
```

$$triad_e := [\,e\_e\_x,\ e\_e\_y,\ e\_e\_z\,]$$

We then transform then multiply by the rotation matrix R_{Euler} to get the triad in the wave propagation frame

```
> triad[w]:=multiply(R[sky],triad[e]);
```

$$\begin{aligned}triad_w := [\,&(\cos(\psi)\sin(\phi)-\sin(\psi)\cos(\theta)\cos(\phi))\,e\_e\_x\\
&+(-\cos(\psi)\cos(\phi)-\sin(\psi)\sin(\phi)\cos(\theta))\,e\_e\_y\\
&+\sin(\psi)\sin(\theta)\,e\_e\_z,\\
&(-\sin(\psi)\sin(\phi)-\cos(\psi)\cos(\theta)\cos(\phi))\,e\_e\_x\\
&+(\sin(\psi)\cos(\phi)-\cos(\psi)\sin(\phi)\cos(\theta))\,e\_e\_y\\
&+\cos(\psi)\sin(\theta)\,e\_e\_z,\\
&-\sin(\theta)\cos(\phi)\,e\_e\_x-\sin(\theta)\sin(\phi)\,e\_e\_y-\cos(\theta)\,e\_e\_z\,]\end{aligned}$$

The individual wave propagation frame vectors can now be written in the earth fixed frame. They are

```
> ind:=[x,y,z]:
  for c1 from 1 to 3 do
  e_w.(ind[c1]):=lin_com(coeff(triad[w][c1],e_e_x),e_ex,coeff(triad[w][
  c1],e_e_y),e_ey,coeff(triad[w][c1],e_e_z),e_ez);
  od; c1:='c1':
```

$$\begin{aligned}e\_wx := \mathrm{table}([\\
&index\_char=[\,\text{-1}\,]\\
&compts=[\,\cos(\psi)\sin(\phi)-\sin(\psi)\cos(\theta)\cos(\phi),\\
&-\cos(\psi)\cos(\phi)-\sin(\psi)\sin(\phi)\cos(\theta),\ \sin(\psi)\sin(\theta)\,]\\
&])\end{aligned}$$

$$\begin{aligned}e\_wy := \mathrm{table}([\\
&index\_char=[\,\text{-1}\,]\\
&compts=[\,-\sin(\psi)\sin(\phi)-\cos(\psi)\cos(\theta)\cos(\phi),\\
&\sin(\psi)\cos(\phi)-\cos(\psi)\sin(\phi)\cos(\theta),\ \cos(\psi)\sin(\theta)\,]\\
&])\end{aligned}$$

$$\begin{aligned}e\_wz := \mathrm{table}([\\
&index\_char=[\,\text{-1}\,]\\
&compts=[\,-\sin(\theta)\cos(\phi),\ -\sin(\theta)\sin(\phi),\ -\cos(\theta)\,]\\
&])\end{aligned}$$

Now we can recalculate the `e_plus` and `e_cross` in the earth fixed frame (results supressed),

```
> `tensor/lin_com/simp`:= proc(x) simplify(x, trig) end:
  e_plus:=lin_com(1,prod(e_wx,e_wx),-1,prod(e_wy,e_wy)):
```

```
  e_cross:=lin_com(1,prod(e_wx,e_wy),1,prod(e_wy,e_wx)):
```

And finally, the gravitational wave tensor in the earth fixed frame (results supressed).

```
> `tensor/lin_com/simp`:= proc(x) collect(simplify(x,
  trig),{h_plus,h_cross}) end:
  h_tensor:=lin_com(h_plus,e_plus,h_cross,e_cross):
```

To verify, calculate the response functions for a detector at the center of the earth with arms along the x_e and y_e axes. Since h_tensor has already been written in the detector frame above, the strain h is just given by $h\_tensor_{1,\ 1} - h\_tensor_{2,\ 2}$. F_plus and F_cross are then simply the coefficients of h_{plus} and h_{cross} respectively.

```
> F_plus:=coeff(get_compts(h_tensor)[1,1]/2-get_compts(h_tensor)[2,2]/2
  ,h_plus);
  F_cross:=coeff(get_compts(h_tensor)[1,1]/2-get_compts(h_tensor)[2,2]/
  2,h_cross);
```

$$F\_plus := \cos(\psi)^2 - 2\cos(\psi)^2\cos(\phi)^2 - 4\cos(\psi)\sin(\phi)\sin(\psi)\cos(\theta)\cos(\phi) + \cos(\theta)^2\cos(\phi)^2 - 2\cos(\theta)^2\cos(\phi)^2\cos(\psi)^2 - \frac{1}{2} + \cos(\phi)^2 - \frac{1}{2}\cos(\theta)^2 + \cos(\theta)^2\cos(\psi)^2$$

$$F\_cross := -\cos(\psi)\sin(\psi) + 2\cos(\psi)\sin(\psi)\cos(\phi)^2 - 4\cos(\psi)^2\sin(\phi)\cos(\theta)\cos(\phi) + 2\cos(\theta)\cos(\phi)\sin(\phi) + 2\sin(\psi)\cos(\theta)^2\cos(\phi)^2\cos(\psi) - \sin(\psi)\cos(\theta)^2\cos(\psi)$$

This is a bit messy, so let’s simplify

```
> tmp1:=expand(cos(2*phi))=cos(2*phi):
  tmp2:=expand(cos(2*psi))=cos(2*psi):
  tmp3:=expand(sin(2*phi))=sin(2*phi):
  tmp4:=expand(sin(2*psi))=sin(2*psi):
  F_plus:=algsubs(tmp1,algsubs(tmp2,algsubs(tmp3,algsubs(tmp4,F_plus)))
  );
  F_cross:=algsubs(tmp1,algsubs(tmp2,algsubs(tmp3,algsubs(tmp4,F_cross)
  )));
```

$$F\_plus := -\cos(\theta)\sin(2\,\psi)\sin(2\,\phi) - \frac{1}{2}\cos(\theta)^2\cos(2\,\psi)\cos(2\,\phi) - \frac{1}{2}\cos(2\,\psi)\cos(2\,\phi)$$

$$F\_cross := \frac{1}{2}\cos(\theta)^2\sin(2\,\psi)\cos(2\,\phi) + \frac{1}{2}\sin(2\,\psi)\cos(2\,\phi) - \cos(\theta)\sin(2\,\phi)\cos(2\,\psi)$$

This is exactly what Anderson, Brady, Creighton and Flanagan get!

Location and Arm Orientation Vectors for Detectors

The primary advantage of the earth fixed frame is that one can conveniently express all the vectors describing the detector arms in it, and these vectors, once determined, are fixed for all time. Our goal, therefore, is to find for each detector three vectors:

- a *location vector* describing the position of the corner station and
- two *arm orientation vectors* describing the orientation of the arms.

These can then be used with the gravitational wave tensor in the earth fixed frame to obtain the strain produced by the gravitational wave anywhere. Information which can be used to construct these vectors for each detector is (mostly) contained within the data frames (not to be confused with coordinate frames) produced by the detector. However, this information is not specified in terms of the earth fixed coordinates. The information used to reconstruct the location vector is given in terms of the earth model WGS-84 coordinates. The information needed for the arm orientation vectors is encoded in a set of arm orientation angles.

Arm orientation angle definitions

There are four arm orientation angles, ψ_1 and ψ_2 which are found in data frames, and ω_1 and ω_2 which are not in the data frames at the present time. We note in passing that while three angles are sufficient to specify the orientation of any frame (and hence of the detector arms), we will follow LIGO-T980044-08-E in using a four-angle system. The four orientation angles most easily defined in the detector frame. They are defined to be:

- ψ_1: the angle North of East of the projection of one arm (arm 1) onto the local horizontal.
- ψ_2: the angle North of East of the projection of the other arm (arm 2) onto the local horizontal.
- ω_1: the angle of smallest magnitude between arm 1 and the local horizontal. Angles above (below) the horizontal are +ve (-ve).
- ω_2: the angle of smallest magnitude between arm 1 and the local horizontal. Angles above (below) the horizontal are +ve (-ve).

We furthermore require the transformation from earth model WGS-84 coordinates to earth fixed coordinates,

Transformation from earth model WGS-84 coordinates to earth fixed coordinates.

This transformation from (l, λ, h), the earth model WGS-84 coordinates to the fixed earth frame coordinates (x_e, y_e, z_e) is described in LIGO-T980044-08-E. The transformation is

```
> M2E_x:=x[e]=(R+h)*cos(l)*cos(lambda);
  M2E_y:=y[e]=(R+h)*cos(l)*sin(lambda);
  M2E_z:=z[e]=(b^2*R/a^2+h)*sin(l);
  earthModel_to_earthFixed:=[M2E_x,M2E_y,M2E_z];
```

$$M2E\_x := x_e = (R+h)\cos(l)\cos(\lambda)$$

$$M2E\_y := y_e = (R+h)\cos(l)\sin(\lambda)$$

$$M2E\_z := z_e = \left(\frac{b^2 R}{a^2} + h\right)\sin(l)$$

$$earthModel\_to\_earthFixed := \left[x_e = (R+h)\cos(l)\cos(\lambda),\ y_e = (R+h)\cos(l)\sin(\lambda),\ z_e = \left(\frac{b^2 R}{a^2} + h\right)\sin(l)\right]$$

where R is the local radius of curvature, defined as

```
> Rdef:=R=a^2/sqrt(a^2*cos(l)^2+b^2*sin(l)^2);
```

$$Rdef := R = \frac{a^2}{\sqrt{a^2 \cos(l)^2 + b^2 \sin(l)^2}}$$

and the detector frame basis vectors expressed in terms of the earth fixed frame basis.

Detector frame basis vectors in the earth fixed frame

The components of the detector frame basis vectors in the earth fixed frame basis are can be calculated in terms of the derivatives of the transformation from earth model WGS-84 coordinates to earth fixed coordinates (i.e. the Jacobian of the transformation). For the λ and h derivatives, which give us the components of e_d_x and e_d_z respectively, this is relatively trivial. The components of e_d_y are considerably more difficult. First, we calculate the components of e_d_x in the earth fixed frame.

λ derivatives and e_d_x

First the derivatives with respect to λ.

```
> M2E_derivs[lambda] := [ diff(rhs(M2E_x),lambda),
  diff(rhs(M2E_y),lambda), diff(rhs(M2E_z),lambda) ];
```

$$M2E\_derivs_{\lambda} := [-(R+h)\cos(l)\sin(\lambda),\ (R+h)\cos(l)\cos(\lambda),\ 0]$$

This gives an un-normalized vector. The normalization factor is:

```
> e_dx_norm:=simplify(sqrt(algsubs(sin(lambda)^2+cos(lambda)^2=1,
  factor(M2E_derivs[lambda][1]^2+
  M2E_derivs[lambda][2]^2+M2E_derivs[lambda][3]^2))),symbolic);
```

$$e\_dx\_norm := \cos(l)\,(R+h)$$

The components of the detector frame x basis vector in the earth fixed basis are therefore:

```
> e_d_x:=array([M2E_derivs[lambda][i]/e_dx_norm$i=1..3]);
```

$$e\_d\_x := [-\sin(\lambda),\ \cos(\lambda),\ 0]$$

Now the components of e_d_z.

h derivatives and e_d_z

Derivatives with respect to h.

```
> M2E_derivs[h] := [ diff(rhs(M2E_x),h), diff(rhs(M2E_y),h),
  diff(rhs(M2E_z),h) ];
```

$$M2E\_derivs_h := [\cos(l)\cos(\lambda), \cos(l)\sin(\lambda), \sin(l)]$$

This is, in fact normalized:

```
> e_dz_norm:=sqrt(simplify(M2E_derivs[h][1]^2+
  M2E_derivs[h][2]^2+M2E_derivs[h][3]^2,trig));
```

$$e\_dz\_norm := 1$$

The components of the detector frame x; basis vector in the earth fixed basis are therefore:

```
> e_d_z:=array([M2E_derivs[h][i]/e_dz_norm$i=1..3]);
```

$$e\_d\_z := [\cos(l)\cos(\lambda), \cos(l)\sin(\lambda), \sin(l)]$$

We could find the components of *e_d_y* by taking a cross product of the other two basis vectors, to complete a right handed frame. However, for pedagogy's sake, we calculate it using derivatives as well.

l derivatives and *e_d_y*

First, let's compute the derivatives (results supressed).

```
> M2E_derivs[l] := [ diff(rhs(subs(Rdef,M2E_x)),l),
  diff(rhs(subs(Rdef,M2E_y)),l), diff(rhs(subs(Rdef,M2E_z)),l) ]:
```

Since we will be normalizing anyway, we can always rescale. When algebraic functions are involved, one will tend to get repeated multiplicative factors. A simple way to remove and identify these is to consider ratios of components. We do so to get simple forms for the *x* and *y* components of *e_d_y*.

```
> Lx:=numer(normal(M2E_derivs[l][1]/M2E_derivs[l][2]));Ly:=denom(
  normal(M2E_derivs[l][1]/M2E_derivs[l][2]));
```

$$Lx := \cos(\lambda)$$

$$Ly := \sin(\lambda)$$

Now remove the same factor from the *z* component.

```
> Lz:=algsubs(sin(l)^2+cos(l)^2=1,normal(M2E_derivs[l][3]/M2E_der
  ivs[l][1])*Lx);
```

$$Lz := -\frac{\cos(l)}{\sin(l)}$$

The normalization factor is

```
> e_dy_norm:=simplify(sqrt(algsubs(sin(l)^2+cos(l)^2=1,algsubs(si
  n(lambda)^2+cos(lambda)^2=1,normal(Lx^2+Ly^2+Lz^2)))),symbolic)
  ;
```

$$e\_dy\_norm := \frac{1}{\sin(l)}$$

Up to a sign (which we may have lost by rescaling by removing a negative multiplicative factor above), the vector is

```
> e_d_y_tmp:=array([Lx/e_dy_norm,Ly/e_dy_norm,Lz/e_dy_norm]);
```

$$e\_d\_y\_tmp := [\cos(\lambda)\sin(l), \sin(\lambda)\sin(l), -\cos(l)]$$

We can check the sign by computing (*e_d_x X e_d_y*) and seeing whether we get *e_d_z* or *−e_d_z*.

```
> crossprod(e_d_x,e_d_y_tmp);
```

$$[-\cos(l)\cos(\lambda), -\cos(l)\sin(\lambda),$$
$$-\sin(\lambda)^2\sin(l)-\cos(\lambda)^2\sin(l)]$$

This is $-\mathit{e\_d\_z}$, so $\mathit{e\_d\_y}$ is actually given by

```
> e_d_y:=array([-e_d_y_tmp[1],-e_d_y_tmp[2],-e_d_y_tmp[3]]);
```

$$\mathit{e\_d\_y} := [-\cos(\lambda)\sin(l), -\sin(\lambda)\sin(l), \cos(l)]$$

Location vector

The location vector Λ for a detector is given simply by the WGS-84 to earth fixed coordinate transformation,

```
> Lambda:=[rhs(M2E_x),rhs(M2E_y),rhs(M2E_z)];
```

$$\Lambda := \Bigg[(R+h)\cos(l)\cos(\lambda), (R+h)\cos(l)\sin(\lambda),$$
$$\left(\frac{b^2 R}{a^2}+h\right)\sin(l)\Bigg]$$

where R is given above, and the parameters a and b are also given above.

Orientation vectors

For each of the two arms, we want to have unit orientation vectors, which will determine the orientation of the arms in terms of the orientation angles. In the detector coordinate frame, the angles $\psi_{1,\ 2}$ play the role of azimuthal angles for arms 1 and 2 respectively, while $\omega_{1,\ 2}$ are the polar complementary angles to the polar angle. Thus, it is straightforward to see that:

```
> Omega_d[1]:=vector([cos(psi1)*cos(omega1),sin(psi1)*cos(omega1),si
  n(omega1)]);
  Omega_d[2]:=vector([cos(psi2)*cos(omega2),sin(psi2)*cos(omega2),si
  n(omega2)]);
```

$$\mathit{Omega\_d}_1 := [\cos(\psi1)\cos(\omega1), \sin(\psi1)\cos(\omega1), \sin(\omega1)]$$

$$\mathit{Omega\_d}_2 := [\cos(\psi2)\cos(\omega2), \sin(\psi2)\cos(\omega2), \sin(\omega2)]$$

```
> transpose(evalm(Omega_d[1]));
```

$$\mathrm{transpose}([\cos(\psi1)\cos(\omega1), \sin(\psi1)\cos(\omega1), \sin(\omega1)])$$

To transform these into the earth fixed frame, we need to use the detector frame to earth fixed frame transformation.

```
> for C1 from 1 to 2 do
     Omega_.(C1) :=
  evalm(Omega_d[C1][1]*e_d_x+Omega_d[C1][2]*e_d_y+Omega_d[C1][3]*e_d
  _z)
  od;C1:='C1':
```

$$\mathit{Omega\_1} := [-\cos(\psi1)\cos(\omega1)\sin(\lambda)$$
$$-\sin(\psi1)\cos(\omega1)\cos(\lambda)\sin(l)+\sin(\omega1)\cos(l)\cos(\lambda),$$

$$\cos(\psi 1)\cos(\omega 1)\cos(\lambda)-\sin(\psi 1)\cos(\omega 1)\sin(\lambda)\sin(l)$$
$$+\sin(\omega 1)\cos(l)\sin(\lambda),$$
$$\sin(\psi 1)\cos(\omega 1)\cos(l)+\sin(\omega 1)\sin(l)\,]$$

$$\mathit{Omega\_2} := [\,-\cos(\psi 2)\cos(\omega 2)\sin(\lambda)$$
$$-\sin(\psi 2)\cos(\omega 2)\cos(\lambda)\sin(l)+\sin(\omega 2)\cos(l)\cos(\lambda),$$
$$\cos(\psi 2)\cos(\omega 2)\cos(\lambda)-\sin(\psi 2)\cos(\omega 2)\sin(\lambda)\sin(l)$$
$$+\sin(\omega 2)\cos(l)\sin(\lambda),$$
$$\sin(\psi 2)\cos(\omega 2)\cos(l)+\sin(\omega 2)\sin(l)\,]$$

These arm direction unit vectors allow us to construct the *response tensor* in the earth-fixed frame. As a matrix, this is one-half the difference between the matrices generated by taking the outer product of each arm vector with itself. (results supressed, but available below under "Location and Orientation vectors for known detectors")

```
> Omega_1m:=convert(Omega_1,matrix):
  Omega_2m:=convert(Omega_2,matrix):
  Omega_1M:=transpose(Omega_1m): Omega_2M:=transpose(Omega_2m):
  response_tensor:=evalm((Omega_1m &* Omega_1M-Omega_2m &*
  Omega_2M)/2):
```

Location and Orientation vectors for known detectors

First, put all vectors into a list.

```
> Det_vecs:=[Lambda,convert(Omega_1,list),convert(Omega_2,list),eval
  (response_tensor)];
  Det_vecs:=subs(WGS84_params, subs(Rdef,Det_vecs)):
```

$$\mathit{Det\_vecs} := \Big[\Big[\,(R+h)\cos(l)\cos(\lambda),\ (R+h)\cos(l)\sin(\lambda),$$
$$\left(\frac{b^2\,R}{a^2}+h\right)\sin(l)\Big],\ [\,-\cos(\psi 1)\cos(\omega 1)\sin(\lambda)$$
$$-\sin(\psi 1)\cos(\omega 1)\cos(\lambda)\sin(l)+\sin(\omega 1)\cos(l)\cos(\lambda),$$
$$\cos(\psi 1)\cos(\omega 1)\cos(\lambda)-\sin(\psi 1)\cos(\omega 1)\sin(\lambda)\sin(l)$$
$$+\sin(\omega 1)\cos(l)\sin(\lambda),$$
$$\sin(\psi 1)\cos(\omega 1)\cos(l)+\sin(\omega 1)\sin(l)\,],\ [$$
$$-\cos(\psi 2)\cos(\omega 2)\sin(\lambda)-\sin(\psi 2)\cos(\omega 2)\cos(\lambda)\sin(l)$$
$$+\sin(\omega 2)\cos(l)\cos(\lambda),\ \cos(\psi 2)\cos(\omega 2)\cos(\lambda)$$
$$-\sin(\psi 2)\cos(\omega 2)\sin(\lambda)\sin(l)+\sin(\omega 2)\cos(l)\sin(\lambda),$$
$$\sin(\psi 2)\cos(\omega 2)\cos(l)+\sin(\omega 2)\sin(l)\,],$$
$$\Big[\frac{1}{2}\,(-\cos(\psi 1)\cos(\omega 1)\sin(\lambda)-\sin(\psi 1)\cos(\omega 1)\cos(\lambda)\sin(l)$$

$$+\sin(\omega1)\cos(l)\cos(\lambda))^{2}-\frac{1}{2}(-\cos(\psi2)\cos(\omega2)\sin(\lambda)$$

$$-\sin(\psi2)\cos(\omega2)\cos(\lambda)\sin(l)+\sin(\omega2)\cos(l)\cos(\lambda))^{2},$$

$$\frac{1}{2}(-\cos(\psi1)\cos(\omega1)\sin(\lambda)-\sin(\psi1)\cos(\omega1)\cos(\lambda)\sin(l)$$

$$+\sin(\omega1)\cos(l)\cos(\lambda))(\cos(\psi1)\cos(\omega1)\cos(\lambda)$$

$$-\sin(\psi1)\cos(\omega1)\sin(\lambda)\sin(l)+\sin(\omega1)\cos(l)\sin(\lambda))-\frac{1}{2}$$

$$(-\cos(\psi2)\cos(\omega2)\sin(\lambda)-\sin(\psi2)\cos(\omega2)\cos(\lambda)\sin(l)$$

$$+\sin(\omega2)\cos(l)\cos(\lambda))(\cos(\psi2)\cos(\omega2)\cos(\lambda)$$

$$-\sin(\psi2)\cos(\omega2)\sin(\lambda)\sin(l)+\sin(\omega2)\cos(l)\sin(\lambda)),\ \frac{1}{2}$$

$$(-\cos(\psi1)\cos(\omega1)\sin(\lambda)-\sin(\psi1)\cos(\omega1)\cos(\lambda)\sin(l)$$

$$+\sin(\omega1)\cos(l)\cos(\lambda))$$

$$(\sin(\psi1)\cos(\omega1)\cos(l)+\sin(\omega1)\sin(l))-\frac{1}{2}\ ($$

$$-\cos(\psi2)\cos(\omega2)\sin(\lambda)-\sin(\psi2)\cos(\omega2)\cos(\lambda)\sin(l)$$

$$+\sin(\omega2)\cos(l)\cos(\lambda))$$

$$(\sin(\psi2)\cos(\omega2)\cos(l)+\sin(\omega2)\sin(l))\Bigg]$$

$$\Bigg[\frac{1}{2}(-\cos(\psi1)\cos(\omega1)\sin(\lambda)-\sin(\psi1)\cos(\omega1)\cos(\lambda)\sin(l)$$

$$+\sin(\omega1)\cos(l)\cos(\lambda))(\cos(\psi1)\cos(\omega1)\cos(\lambda)$$

$$-\sin(\psi1)\cos(\omega1)\sin(\lambda)\sin(l)+\sin(\omega1)\cos(l)\sin(\lambda))-\frac{1}{2}$$

$$(-\cos(\psi2)\cos(\omega2)\sin(\lambda)-\sin(\psi2)\cos(\omega2)\cos(\lambda)\sin(l)$$

$$+\sin(\omega2)\cos(l)\cos(\lambda))(\cos(\psi2)\cos(\omega2)\cos(\lambda)$$

$$-\sin(\psi2)\cos(\omega2)\sin(\lambda)\sin(l)+\sin(\omega2)\cos(l)\sin(\lambda)),\ \frac{1}{2}$$

$$(\cos(\psi1)\cos(\omega1)\cos(\lambda)-\sin(\psi1)\cos(\omega1)\sin(\lambda)\sin(l)$$

$$+\sin(\omega1)\cos(l)\sin(\lambda))^{2}-\frac{1}{2}(\cos(\psi2)\cos(\omega2)\cos(\lambda)$$

$$-\sin(\psi2)\cos(\omega2)\sin(\lambda)\sin(l)+\sin(\omega2)\cos(l)\sin(\lambda))^{2},$$

$$\frac{1}{2}(\cos(\psi1)\cos(\omega1)\cos(\lambda)-\sin(\psi1)\cos(\omega1)\sin(\lambda)\sin(l)$$

$$+\sin(\omega 1)\cos(l)\sin(\lambda))$$

$$(\sin(\psi 1)\cos(\omega 1)\cos(l)+\sin(\omega 1)\sin(l))-\frac{1}{2}\,($$

$$\cos(\psi 2)\cos(\omega 2)\cos(\lambda)-\sin(\psi 2)\cos(\omega 2)\sin(\lambda)\sin(l)$$

$$+\sin(\omega 2)\cos(l)\sin(\lambda))$$

$$(\sin(\psi 2)\cos(\omega 2)\cos(l)+\sin(\omega 2)\sin(l))\,\Big]$$

$$\Big[\frac{1}{2}\,(-\cos(\psi 1)\cos(\omega 1)\sin(\lambda)-\sin(\psi 1)\cos(\omega 1)\cos(\lambda)\sin(l)$$

$$+\sin(\omega 1)\cos(l)\cos(\lambda))$$

$$(\sin(\psi 1)\cos(\omega 1)\cos(l)+\sin(\omega 1)\sin(l))-\frac{1}{2}\,($$

$$-\cos(\psi 2)\cos(\omega 2)\sin(\lambda)-\sin(\psi 2)\cos(\omega 2)\cos(\lambda)\sin(l)$$

$$+\sin(\omega 2)\cos(l)\cos(\lambda))$$

$$(\sin(\psi 2)\cos(\omega 2)\cos(l)+\sin(\omega 2)\sin(l))\,,\;\frac{1}{2}\,($$

$$\cos(\psi 1)\cos(\omega 1)\cos(\lambda)-\sin(\psi 1)\cos(\omega 1)\sin(\lambda)\sin(l)$$

$$+\sin(\omega 1)\cos(l)\sin(\lambda))$$

$$(\sin(\psi 1)\cos(\omega 1)\cos(l)+\sin(\omega 1)\sin(l))-\frac{1}{2}\,($$

$$\cos(\psi 2)\cos(\omega 2)\cos(\lambda)-\sin(\psi 2)\cos(\omega 2)\sin(\lambda)\sin(l)$$

$$+\sin(\omega 2)\cos(l)\sin(\lambda))$$

$$(\sin(\psi 2)\cos(\omega 2)\cos(l)+\sin(\omega 2)\sin(l))\,,$$

$$\frac{1}{2}\,(\sin(\psi 1)\cos(\omega 1)\cos(l)+\sin(\omega 1)\sin(l))^2$$

$$-\frac{1}{2}\,(\sin(\psi 2)\cos(\omega 2)\cos(l)+\sin(\omega 2)\sin(l))^2\,\Big]\Big]$$

Now the detectors:

Reference to make sure we're getting it right (REF)

Let's try a detector with arms due east and due south at 0 degrees longitude and 0 degrees latitude ... should be `[ [ a, 0, 0 ], [ 0, 1, 0 ], [ 0, 0, 1 ] ]`.

```
> params_REF:=[
      l          = 0,
      lambda     = 0,
      h          = 0,
      psi1       = 0,
      psi2     = Pi/2,
```

```
      omega1   = 0,
      omega2   = 0
   ];
```

```
   params_REF := [ l = 0,  λ = 0,  h = 0,  ψ1 = 0,  ψ2 = 1/2 π,  ω1 = 0,  ω2 = 0 ]
```

```
> DV_REF:=evalf(subs(params_REF,Det_vecs),12);
```

```
DV_REF := [ [ .637813700002 10^7, 0, 0 ], [ 0, 1., 0 ],

    [ 0, -.510338076868 10^-11, 1. ],
    [0              0                          0          ]
    [0        .500000000000       .255169038434 10^-11    ]
    [0    .255169038434 10^-11      -.500000000000        ] ]
```

LIGO-Hanford Observatory (LHO)

Reference:
William Althouse, Larry Jones, Albert Lazzarini (1999)
"Determination of Global and Local Coordinate Axes for the LIGO Sites"
LIGO-T980044-08-E

```
> params_LHO:=[
      h          = 142.554,
      l          = (46 + 27/60 + 18.528/3600)*Pi/180,
      lambda     = (240 + 35/60 + 32.4343/3600)*Pi/180,
      psi1     = 125.9994*Pi/180,
      psi2     = 215.9994*Pi/180,
      omega1   = -6.195e-4,
      omega2   = 1.25e-5
   ];
```

```
params_LHO := [ h = 142.554,  l = .2580841482 π,  λ = 1.336624127 π,

    ψ1 = .6999966667 π,  ψ2 = 1.199996667 π,  ω1 = -.0006195,

    ω2 = .0000125 ]
```

```
> DV_LHO:=evalf(subs(params_LHO,Det_vecs),12);
```

```
DV_LHO := [

    [ -.216141492611 10^7, -.383469517823 10^7, .460035022722 10^7 ],
    [ -.223892661573, .799830627548, .556904878171 ],
    [ -.913978185537, .0260940390132, -.404923421764 ],
```

```
[ -.392614099866     -.0776134127892     -.247389044862 ]
[-.0776134127892      .319524066946      .227997832878  ]
[ -.247389044862      .227997832878      .0730900329195 ]
```

LIGO-Livingston Observatory (LLO)

Reference:
William Althouse, Larry Jones, Albert Lazzarini (1999)
"Determination of Global and Local Coordinate Axes for the LIGO Sites"
LIGO-T980044-08-E

```
> params_LLO:=[
     h          = -6.574,
     l          = (30 + 33/60 + 46.4196/3600)*Pi/180,
     lambda     = (269 + 13/60 + 32.7346/3600)*Pi/180,
     psi1     = 197.7165*Pi/180,
     psi2     = 287.7165*Pi/180,
     omega1   = -3.121e-4,
     omega2   = -6.107e-4
  ];
params_LLO := [ h = -6.574, l = .1697938574 π, λ = 1.495698664 π,
    ψ1 = 1.098425000 π, ψ2 = 1.598425000 π, ω1 = -.0003121,
    ω2 = -.0006107 ]
```

```
> DV_LLO:=evalf(subs(params_LLO,Det_vecs),12);
DV_LLO := [

    [ -74276.0534461, -.549628371972 10^7, .322425701722 10^7 ],
    [ -.954574121757, -.141580771885, -.262189113245 ],
    [ .297741568158, -.487910336926, -.820544612871 ],
    [.411280856259     .140210264909      .247294591145 ]
    [.140210264909     -.109005690956     -.181615630748]
    [.247294591145     -.181615630748     -.302275165304] ]
```

VIRGO Interferometer (VIRGO)

Reference:
Benoit Mours, E-mail: mours@lapp.in2p3.fr

```
> params_VIRGO:=[
     h         = 51.884,
     l         = (43 + 37/60 + 53.0921/3600)*Pi/180,
     lambda    = (10 + 30/60 + 16.1878/3600)*Pi/180,
     psi1    = 70.5674*Pi/180,
     psi2    = 160.5674*Pi/180,
     omega1 = 0,
     omega2 = 0
  ];
```

```
params_VIRGO := [ h = 51.884,  l = .2423967471 π,  λ = .05835831450 π,
    ψ1 = .3920411111 π,  ψ2 = .8920411111 π,  ω1 = 0,  ω2 = 0 ]
> DV_VIRGO:=evalf(subs(params_VIRGO,Det_vecs),12);
DV_VIRGO  := [

    [ .454637409863 10^7, 842989.697467, .437857696275 10^7 ],
    [ -.700458214851, .208489486220, .682561662696 ],
    [ -.053792553751, -.969081805482, .240804517079 ],
    [ .243874035956      -.0990837792218     -.232576216925 ]
    [-.0990837792218     -.447825839926      .187833103275  ]
    [ -.232576216925      .187833103275      .203951803968  ]]
```

GEO-600 Interferometer (GEO600)

Reference:
http://www.geo600.uni-hannover.de/geo600/project/location.html

```
> params_GEO600:=[
      h         = 114.425,
      l         = (52 + 14/60 + 42.528/3600)*Pi/180,
      lambda    = (9 + 48/60 + 25.894/3600)*Pi/180,
      psi1    = 21.6117*Pi/180,
      psi2    = 115.9431*Pi/180,
      omega1 = 0,
      omega2 = 0
  ];
params_GEO600 := [ h = 114.425,  l = .2902508148 π,
    λ = .05448440432 π,  ψ1 = .1200650000 π,  ψ2 = .6441283333 π,  ω1 = 0,
    ω2 = 0 ]
> DV_GEO600:=evalf(subs(params_GEO600,Det_vecs),12);
DV_GEO600  := [

    [ .385630994953 10^7, 666598.956352, .501964141692 10^7 ],
    [ -.445306769034, .866513541301, .225513113130 ],
    [ -.626057567788, -.552186095145, .550583724918 ],
    [-.0968249798190     -.365782314548      .122137295952 ]
    [ -.365782314548      .222968116793      .249717421691 ]
    [ .122137295952      .249717421691     -.126143136976 ]]
```

TAMA-300 Interferometer (TAMA300)

Reference:
Masa-Katsu Fujimoto (1995) (unpublished), E-mail: fujimoto@gravity.mtk.nao.ac.jp

```
> params_TAMA300:=[
      h         = 90,
      l         = (35 + 40/60 + 35.6/3600)*Pi/180,
      lambda    = (139 + 32/60 + 9.8/3600)*Pi/180,
      psi1    = 180*Pi/180,
      psi2    = 270*Pi/180,
      omega1 = 0,
      omega2 = 0
  ];
```

$$params\_TAMA300 := \left[h = 90,\ l = .1982030864\ \pi,\ \lambda = .7752003083\ \pi,\ \psi 1 = \pi,\ \psi 2 = \frac{3}{2}\ \pi,\ \omega 1 = 0,\ \omega 2 = 0 \right]$$

```
> DV_TAMA300:=evalf(subs(params_TAMA300,Det_vecs),12);
```

$$DV\_TAMA300 := \Big[[-.394640898771\ 10^{7}, .336625903242\ 10^{7}, .369915069189\ 10^{7}],$$
$$[.648969406115, .760814504283, -.167956993920\ 10^{-12}],$$
$$[-.443713768767, .378484715231, -.812322233933],$$
$$\left[\begin{array}{ccc} .112139690741 & .330842108212 & -.180219279936 \\ .330842108212 & .217794015133 & .153725774693 \\ -.180219279936 & .153725774693 & -.329933705871 \end{array}\right]\Big]$$

Caltech-40 Interferometer (CIT40)

Reference:
B. Allen, "Gravitational Wave Detector Sites," gr-qc/9607075 (1996).

```
> params_CIT40:=[
      h         = 0,
      l         = 34.17*Pi/180,
      lambda    = -118.13*Pi/180,
      psi1    = 270*Pi/180,
      psi2    = 0*Pi/180,
      omega1 = 0,
      omega2 = 0
  ];
```

$$params\_CIT40 := \left[h = 0,\ l = .1898333333\ \pi,\ \lambda = -.6562777778\ \pi,\ \psi 1 = \frac{3}{2}\ \pi,\ \psi 2 = 0,\ \omega 1 = 0,\ \omega 2 = 0 \right]$$

```
> DV_CIT40:=evalf(subs(params_CIT40,Det_vecs),12);
```

```
DV_CIT40 := [

   [ -.249064958399 10^7, -.465869968229 10^7, .356206411337 10^7 ],
   [ -.264803316317, -.495308185286, -.827374767114 ],
   [ .881880123832, -.471473697241, 0 ],
   [-.353795878239     .273471266285     .109545791085]
   [ .273471266285     .011521375611     .204902747226]
   [ .109545791085     .204902747226     .342274502629]]
```

MPQ-30 (MPQ30)

Reference:
B. Allen, "Gravitational Wave Detector Sites," gr-qc/9607075 (1996).

```
> params_MPQ30:=[
      h         = 0,
      l         = 35.57*Pi/180,
      lambda    = 139.47*Pi/180,
      psi1   = 132*Pi/180,
      psi2   = 225*Pi/180,
      omega1 = 0,
      omega2 = 0
  ];
```

$$params\_MPQ30 := \left[h = 0,\ l = .1976111111\ \pi,\ \lambda = .7748333333\ \pi,\ \psi 1 = \frac{11}{15}\pi,\ \psi 2 = \frac{5}{4}\pi,\ \omega 1 = 0,\ \omega 2 = 0 \right]$$

```
> DV_MPQ30:=evalf(subs(params_MPQ30,Det_vecs),12);
```

```
DV_MPQ30 := [

   [ -.394770376318 10^7, .337523393974 10^7, .368948807262 10^7 ],
   [ .763397142331, .227664424303, .604478050046 ],
   [ .146878798953, .804743690014, -.575164508468 ],
   [ .280600907669      .0277992921345     .272968144106]
   [.0277992921345     -.297890658262      .300239078089]
   [ .272968144106      .300239078089      .017289750593]]
```

TAMA-20 (TAMA20)

Reference:
B. Allen, "Gravitational Wave Detector Sites," gr-qc/9607075 (1996).

```
> params_TAMA20:=[
      h         = 0,
```

```
        l          = 35.68*Pi/180,
        lambda     = 139.54*Pi/180,
        psi2    = 45*Pi/180,
        psi1    = 135*Pi/180,
        omega1 = 0,
        omega2 = 0
    ];
```

$$params\_TAMA20 := \left[h = 0,\ l = .1982222222\ \pi,\ \lambda = .7752222222\ \pi,\ \psi2 = \frac{1}{4}\ \pi,\ \psi1 = \frac{3}{4}\ \pi,\ \omega1 = 0,\ \omega2 = 0\right]$$

```
> DV_TAMA20:=evalf(subs(params_TAMA20,Det_vecs),12);
```

$$DV\_TAMA20 := \left[[-.394641550942\ 10^{7},\ .336579525359\ 10^{7},\ .369940864286\ 10^{7}],\ [.772651341534,\ .270378770836,\ .574373767425],\ [-.145055918411,\ -.805638601252,\ .574373767424],\ \begin{bmatrix} .287974438054 & .0460229363990 & .263553488145 \\ .0460229363990 & -.287974438055 & .309018075911 \\ .263553488145 & .309018075911 & 0 \end{bmatrix} \right]$$

Glasgow-10 (1995) (G1095)

Reference:

B. Allen, "Gravitational Wave Detector Sites," gr-qc/9607075 (1996).

```
> params_G1095:=[
        h          = 0,
        l          = 55.87*Pi/180,
        lambda     = -4.28*Pi/180,
        psi1    = 152*Pi/180,
        psi2    = 242*Pi/180,
        omega1 = 0,
        omega2 = 0
    ];
```

$$params\_G1095 := \left[h = 0,\ l = .3103888889\ \pi,\ \lambda = -.02377777778\ \pi,\ \psi1 = \frac{38}{45}\ \pi,\ \psi2 = \frac{121}{90}\ \pi,\ \omega1 = 0,\ \omega2 = 0\right]$$

```
> DV_G1095:=evalf(subs(params_G1095,Det_vecs),12);
```

$$DV\_G1095 := \Bigg[$$

```
                   7                                   7
[ .357683021395 10 , -267687.639613, .525633471276 10  ],

[ -.453424109740, -.851482838592, .263407578267 ],

[ .693799397147, -.522707959431, -.495397603614 ],

[-.137882090094     .374368657593      .112135606031  ]]
[ .374368657593     .225899706782     -.241617651471  ]]
[ .112135606031    -.241617651471    -.0880176166887  ]]
```

Response Functions and a Simple Example

We now have the information necessary in to calculate the beam pattern functions for a gravitational wave that arrives from a source at known declination and right ascension at a given GMST and with a given polarization angle ψ. We begin by defining a procedure that will return the required results as a list.

Definition of the "Response" procedure

MAPLE Response() procedure. '#' denotes the beginning of a comment. Comments end at carriage returns.

```
> Response:= proc(
     dec,         # declination
     ra,          # right ascension
     time,        # Greenwich mean sidereal time
     pol_angle,   # polarization angle
     lat,         # North latitude
     long,        # East longitude
     height,      # height above reference ellipsoid
     psi_arm1,    # arm 1 polarization angle
     psi_arm2,    # arm 2 polarization angle
     omega_arm1,  # arm 1 angle to horizontal
     omega_arm2)  # arm 2 angle to horizontal

  # specify the local and global variables. Note that delta, alpha
  etc. MUST be global.
     local parms,arm_vec1,arm_vec2,strain,F_p,F_c;
     global delta,alpha,GMST,psi,l,lambda,h,

  WGS84_params,Rdef,sky_to_celestial,Omega,h_tensor,R,F_plus,F_cross
  ;

  # create a set of parameter values to with which to evaluate.
     parms:={
        delta=dec,
        alpha=ra,
        GMST=time,
        psi=pol_angle,
        l=lat,
        lambda=long,
        h=height,
```

```
      psi1=psi_arm1,
      psi2=psi_arm2,
      omega1=omega_arm1,
      omega2=omega_arm2
   };

# create the two arm vectors
   arm_vec1:=create([1], subs(WGS84_params, subs(Rdef, subs(parms,
op(Omega_1)))));
   arm_vec2:=create([1], subs(WGS84_params, subs(Rdef, subs(parms,
op(Omega_2)))));

# dot the strain vectors into h_tensor to get the strain in the
detector
   strain:= get_compts( prod(h_tensor,

lin_com(1,prod(arm_vec1,arm_vec1),-1,prod(arm_vec2,arm_vec2)),
[1,1], [2,2]));

# extract response functions, i.e. the coefficients of the plus
and cross polarizations
   F_p:=subs(parms,subs(sky_to_celestial,coeff(strain,h_plus)))/2;

F_c:=subs(parms,subs(sky_to_celestial,coeff(strain,h_cross)))/2;

# make sure the labels we are using for the response functions are
not assigned.
   F_plus:= 'F_plus';
   F_cross:='F_cross';

# return a list with the response functions
   RETURN([F_plus=simplify(F_p),F_cross=simplify(F_c)]);
end:
```

Response for reference detector

As a check, calculate the known response functions for a refernce detector and source. The detector is at the equator $l = 0$ and prime meridian $\lambda = 0$, with arms running due North $\psi1 = 0$ and East $\psi2 = \frac{\pi}{2}$, on the reference ellipsoid $h = 0$, and in the local horizontal $[\ \omega1 = 0,\ \omega2 = 0\ ]$. The source is at declination $\delta = 0$ and right ascension $\alpha = 0$, and the wave arrives at the reference detector at $GMST = 0$ and with polarization angle $\psi = 0$. The response functions should be $[\ F\_plus = 1,\ F\_cross = 0\ ]$.

```
> Response_REF:=eval(subs(params_REF,Response(0,0,0,0,l,lambda,h,psi
1,psi2,omega1,omega2)));
```

$$Response\_REF := [\ F\_plus = 1,\ F\_cross = 0\ ]$$

Response Functions for LIGO Hanford (LHO) and LIGO

Livingston (LLO) Detectors

Response functions for Hanford and Livingston detectors for an generic source.

Response for LHO

```
> Response_LHO:=simplify(subs(params_LHO,Response(delta,alpha,GMST,psi,l,lambda,h,psi1,psi2,omega1,omega2)));
```

$$
\begin{aligned}
\mathit{Response\_LHO} := [\, & \mathit{F\_plus} = .3104536563 \cos(\psi) \sin(\psi) \sin(\delta) \\
& - .1552268269 \sin(-1.\ \alpha + \mathit{GMST}) \cos(-1.\ \alpha + \mathit{GMST}) \cos(\delta)^2 \\
& - .6209073106 \cos(\psi)^2 \sin(-1.\ \alpha + \mathit{GMST}) \cos(-1.\ \alpha + \mathit{GMST}) \\
& + .4947780908 \cos(\delta) \sin(\delta) \cos(-1.\ \alpha + \mathit{GMST}) \\
& - .9895561806 \cos(\delta) \sin(\delta) \cos(-1.\ \alpha + \mathit{GMST}) \cos(\psi)^2 \\
& + .4559956647 \cos(\delta) \sin(-1.\ \alpha + \mathit{GMST}) \sin(\delta) \\
& - 1.424276333 \cos(\psi)^2 \cos(-1.\ \alpha + \mathit{GMST})^2 \cos(\delta)^2 \\
& - 1.424276335 \cos(\psi)^2 + .4928680670 \cos(\psi)^2 \cos(\delta)^2 \\
& + .9895561807 \sin(\psi) \cos(\delta) \cos(\psi) \sin(-1.\ \alpha + \mathit{GMST}) \\
& - .6209073102 \sin(\psi) \sin(\delta) \cos(-1.\ \alpha + \mathit{GMST})^2 \cos(\psi) \\
& + .3104536568 \sin(-1.\ \alpha + \mathit{GMST}) \cos(-1.\ \alpha + \mathit{GMST}) \\
& + 2.848552664 \cos(\psi)^2 \cos(-1.\ \alpha + \mathit{GMST})^2 - 2.848552664 \\
& \cos(\psi) \sin(-1.\ \alpha + \mathit{GMST}) \sin(\psi) \sin(\delta) \cos(-1.\ \alpha + \mathit{GMST}) \\
& + .7121381649 \cos(-1.\ \alpha + \mathit{GMST})^2 \cos(\delta)^2 + .3104536565 \\
& \cos(\psi)^2 \sin(-1.\ \alpha + \mathit{GMST}) \cos(-1.\ \alpha + \mathit{GMST}) \cos(\delta)^2 + .7121381662 \\
& - .9119913309 \sin(\psi) \cos(\delta) \cos(\psi) \cos(-1.\ \alpha + \mathit{GMST}) \\
& - 1.424276333 \cos(-1.\ \alpha + \mathit{GMST})^2 - .2464340348 \cos(\delta)^2 \\
& - .9119913309 \cos(\delta) \sin(-1.\ \alpha + \mathit{GMST}) \sin(\delta) \cos(\psi)^2,\ \mathit{F\_cross} \\
& = -.4947780908 \cos(\delta) \sin(-1.\ \alpha + \mathit{GMST}) - .3104536569 \\
& \cos(\psi) \sin(-1.\ \alpha + \mathit{GMST}) \sin(\psi) \cos(-1.\ \alpha + \mathit{GMST}) \cos(\delta)^2 \\
& + 1.424276334 \cos(\psi) \sin(\psi) \\
& + .4559956650 \cos(\delta) \cos(-1.\ \alpha + \mathit{GMST}) \\
& - .4928680678 \cos(\psi) \sin(\psi) \cos(\delta)^2 - .1552268269 \sin(\delta) \\
& + .3104536564 \cos(\psi)^2 \sin(\delta) \\
& - 2.848552664 \cos(\psi)^2 \sin(-1.\ \alpha + \mathit{GMST}) \sin(\delta) \cos(-1.\ \alpha + \mathit{GMST})
\end{aligned}
$$

$$+.6209073092 \cos(\psi) \sin(-1.\ \alpha+GMST) \sin(\psi) \cos(-1.\ \alpha+GMST)$$
$$+1.424276334 \cos(\psi) \sin(\psi) \cos(-1.\ \alpha+GMST)^2 \cos(\delta)^2$$
$$+1.424276334 \sin(\delta) \cos(-1.\ \alpha+GMST) \sin(-1.\ \alpha+GMST)$$
$$+.9895561806 \cos(\psi) \cos(\delta) \sin(\psi) \sin(\delta) \cos(-1.\ \alpha+GMST)$$
$$-.6209073097 \cos(\psi)^2 \sin(\delta) \cos(-1.\ \alpha+GMST)^2$$
$$+.9119913300 \cos(\psi) \cos(\delta) \sin(\psi) \sin(-1.\ \alpha+GMST) \sin(\delta)$$
$$+.3104536564 \sin(\delta) \cos(-1.\ \alpha+GMST)^2$$
$$-2.848552664 \cos(\psi) \sin(\psi) \cos(-1.\ \alpha+GMST)^2$$
$$+.9895561807 \cos(\psi)^2 \cos(\delta) \sin(-1.\ \alpha+GMST)$$
$$-.9119913310 \cos(\psi)^2 \cos(\delta) \cos(-1.\ \alpha+GMST)\ ]$$

Response for LLO

```
> Response_LLO:=simplify(subs(params_LLO,Response(delta,alpha,GMST,p
  si,l,lambda,h,psi1,psi2,omega1,omega2)));
```

$$Response\_LLO := [\ F\_plus =$$
$$1.040573095 \cos(\psi)^2 \cos(-1.\ \alpha+GMST)^2 \cos(\delta)^2$$
$$-.5608410611 \cos(\psi) \sin(\psi) \sin(\delta)$$
$$-.3632312613 \cos(\delta) \sin(-1.\ \alpha+GMST) \sin(\delta)$$
$$-.4945891833 \cos(\delta) \sin(\delta) \cos(-1.\ \alpha+GMST)$$
$$+1.121682122 \cos(\psi)^2 \sin(-1.\ \alpha+GMST) \cos(-1.\ \alpha+GMST)$$
$$+.2804205304 \sin(-1.\ \alpha+GMST) \cos(-1.\ \alpha+GMST) \cos(\delta)^2$$
$$+1.040573095 \cos(\psi)^2 + .3865389495 \cos(\psi)^2 \cos(\delta)^2$$
$$-.5608410613 \sin(-1.\ \alpha+GMST) \cos(-1.\ \alpha+GMST) - .5202865474$$
$$+1.121682123 \sin(\psi) \sin(\delta) \cos(-1.\ \alpha+GMST)^2 \cos(\psi)$$
$$+1.040573094 \cos(-1.\ \alpha+GMST)^2$$
$$-2.081146184 \cos(\psi)^2 \cos(-1.\ \alpha+GMST)^2$$
$$-.5202865471 \cos(-1.\ \alpha+GMST)^2 \cos(\delta)^2 - .1932694746 \cos(\delta)^2 -$$
$$.5608410611 \cos(\psi)^2 \sin(-1.\ \alpha+GMST) \cos(-1.\ \alpha+GMST) \cos(\delta)^2$$
$$+2.081146185$$
$$\cos(\psi) \sin(-1.\ \alpha+GMST) \sin(\psi) \sin(\delta) \cos(-1.\ \alpha+GMST)$$
$$+.7264625220 \sin(\psi) \cos(\delta) \cos(\psi) \cos(-1.\ \alpha+GMST)$$

$$\begin{array}{l}
+.9891783665\,\cos(\delta)\,\sin(\delta)\,\cos(-1.\;\alpha+\mathit{GMST})\,\cos(\psi)^2\\
+.7264625219\,\cos(\delta)\,\sin(-1.\;\alpha+\mathit{GMST})\,\sin(\delta)\,\cos(\psi)^2\\
-.9891783675\,\sin(\psi)\,\cos(\delta)\,\cos(\psi)\,\sin(-1.\;\alpha+\mathit{GMST}),\ \mathit{F\_cross}=\\
2.081146184\,\cos(\psi)^2\,\sin(-1.\;\alpha+\mathit{GMST})\,\sin(\delta)\,\cos(-1.\;\alpha+\mathit{GMST})\\
+1.121682122\,\cos(\psi)^2\,\sin(\delta)\,\cos(-1.\;\alpha+\mathit{GMST})^2\\
-.9891783659\,\cos(\psi)^2\,\cos(\delta)\,\sin(-1.\;\alpha+\mathit{GMST})\\
-.7264625222\,\cos(\psi)\,\cos(\delta)\,\sin(\psi)\,\sin(-1.\;\alpha+\mathit{GMST})\,\sin(\delta)\\
+.4945891833\,\cos(\delta)\,\sin(-1.\;\alpha+\mathit{GMST})\\
+.7264625219\,\cos(\psi)^2\,\cos(\delta)\,\cos(-1.\;\alpha+\mathit{GMST})\\
-.9891783664\,\cos(\psi)\,\cos(\delta)\,\sin(\psi)\,\sin(\delta)\,\cos(-1.\;\alpha+\mathit{GMST})\\
-.3632312613\,\cos(\delta)\,\cos(-1.\;\alpha+\mathit{GMST})\\
-1.040573095\,\cos(\psi)\,\sin(\psi)\\
-.3865389495\,\cos(\psi)\,\sin(\psi)\,\cos(\delta)^2+.2804205304\,\sin(\delta)\\
-.5608410613\,\cos(\psi)^2\,\sin(\delta)\\
-1.040573096\,\sin(\delta)\,\cos(-1.\;\alpha+\mathit{GMST})\,\sin(-1.\;\alpha+\mathit{GMST})\\
+2.081146185\,\cos(\psi)\,\sin(\psi)\,\cos(-1.\;\alpha+\mathit{GMST})^2\\
-1.121682123\,\cos(\psi)\,\sin(-1.\;\alpha+\mathit{GMST})\,\sin(\psi)\,\cos(-1.\;\alpha+\mathit{GMST})\\
-.5608410612\,\sin(\delta)\,\cos(-1.\;\alpha+\mathit{GMST})^2+.5608410612\\
\cos(\psi)\,\sin(-1.\;\alpha+\mathit{GMST})\,\sin(\psi)\,\cos(-1.\;\alpha+\mathit{GMST})\,\cos(\delta)^2\\
-1.040573097\,\cos(\psi)\,\sin(\psi)\,\cos(-1.\;\alpha+\mathit{GMST})^2\,\cos(\delta)^2\,]
\end{array}$$

Times of Arrival

The time of arrival of a gravitational wave at a detector can be calculated if one knows the sky angles of the source, the detectors location vector of the detector, and the GMST at which the wave arrives at the origin of the earth fixed coordinates. The method is straightforward. One uses the sky angles, which are the azimuthal and polar angle for a unit vector pointing to the source, to construct just such a vector. The difference between the distance the wave travels to the coordinate origin and to the detector is thus given by the projection of the detectors location vector onto this unit vector. Finally, the difference in time of arrival is just the difference in distance divided by the speed of light.

The source position unit vector

Because the sky angles $\{\phi,\ \theta\}$ are just the azimuthal and polar angles for this vector, the unit vector is

```
> e_sky:=vector([cos(phi)*sin(theta),sin(phi)*sin(theta),cos(theta)]
  );
```

$$e\_sky := [\sin(\theta)\cos(\phi), \sin(\theta)\sin(\phi), \cos(\theta)]$$

The time delay from the origin to the detector

First, calculate the required inner product

```
> det_orig_dist:=simplify(subs(sky_to_celestial,subs(WGS84_params,su
  bs(Rdef,dotprod(e_sky,convert(Lambda,array),orthogonal)))));
```

$$\begin{aligned}
det\_orig\_dist := (&\\
&.4068063159\ 10^{12} \cos(\delta)\cos(-1.\ \alpha + GMST)\cos(l)\cos(\lambda) +\\
&\cos(\delta)\cos(-1.\ \alpha + GMST)\cos(l)\cos(\lambda)\, h\\
&\sqrt{.27233161\ 10^{8}\cos(l)^{2} + .4040829998\ 10^{10}}\\
&- .4068063159\ 10^{12}\cos(\delta)\sin(-1.\ \alpha + GMST)\cos(l)\sin(\lambda) - 1.\\
&\cos(\delta)\sin(-1.\ \alpha + GMST)\cos(l)\sin(\lambda)\, h\\
&\sqrt{.27233161\ 10^{8}\cos(l)^{2} + .4040829998\ 10^{10}}\\
&+ .4040829998\ 10^{12}\sin(\delta)\sin(l)\\
&+ \sin(\delta)\sin(l)\, h\sqrt{.27233161\ 10^{8}\cos(l)^{2} + .4040829998\ 10^{10}}\ )\\
&\Big/\sqrt{.27233161\ 10^{8}\cos(l)^{2} + .4040829998\ 10^{10}}
\end{aligned}$$

The GMST appearing in the above formula is actually the GMST at which the wave arrives at the detector. However, we lose only higher order corrections by taking it to be the GMST at which the wave arrives at the origin of the earth fixed coordinates. We therefore have that the time difference between the arrival of a wave at the origin and at the detector is

```
> Delta_GMST:= -simplify(evalf(det_orig_dist/(3*10^8)));
```

$$\begin{aligned}
Delta\_GMST := -.1000000000\ 10^{-17}\ (&\\
&.1356021053\ 10^{22}\cos(\delta)\cos(-1.\ \alpha + GMST)\cos(l)\cos(\lambda) +\\
&.3333333333\ 10^{10}\cos(\delta)\cos(-1.\ \alpha + GMST)\cos(l)\cos(\lambda)\, h\\
&\sqrt{.27233161\ 10^{8}\cos(l)^{2} + .4040829998\ 10^{10}}\\
&- .1356021053\ 10^{22}\cos(\delta)\sin(-1.\ \alpha + GMST)\cos(l)\sin(\lambda) -\\
&.3333333333\ 10^{10}\cos(\delta)\sin(-1.\ \alpha + GMST)\cos(l)\sin(\lambda)\, h\\
&\sqrt{.27233161\ 10^{8}\cos(l)^{2} + .4040829998\ 10^{10}}\\
&+ .1346943333\ 10^{22}\sin(\delta)\sin(l) + .3333333333\ 10^{10}\\
&\sin(\delta)\sin(l)\, h\sqrt{.27233161\ 10^{8}\cos(l)^{2} + .4040829998\ 10^{10}}\ )\ \Big/
\end{aligned}$$

$$\sqrt{.27233161\ 10^{8}\ \cos(\,l\,)^{2} + .4040829998\ 10^{10}}$$

An example: Time difference between LIGO Hanford (LHO) and LIGO Livingston (LLO) detectors

Using this result, we find that the time difference between a wave arriving at the Hanford and Livingston observatories for a generic source is

```
> GMST[LHO-LLO]:=evalf(subs(params_LHO,Delta_GMST)-subs(params_LLO,Delta_GMST));
```

$$GMST_{LHO-LLO} := .006957129568\ \cos(\,\delta\,)\ \cos(\,-1.\ \alpha + GMST\,) + .00553862846\ \cos(\,\delta\,)\ \sin(\,-1.\ \alpha + GMST\,) - .00458697738\ \sin(\,\delta\,)$$

References

Clifford M. Will and Alan G. Wiseman, **Physical Review** *D54*, page 4813 (1996).

William Althouse, Larry Jones, Albert Lazzarini, *Determination of Global and Local Coordinate Axes for the LIGO Sites*,
LIGO-T980044-08-E (1999).

Warren G. Anderson, Patrick R. Brady, Jolien D. E. Creighton, and Éanna É. Flanagan, *An excess power statistic for detection of burst sources of gravitational radiation*, gr-qc/0008066